EOS and Neutrino Interactions for Nucleosynthesys

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### EOS

### Density Dependence of EOS

- Pressure of neutron matter pushes neutrons out against surface tension ==> R<sub>n</sub>-R<sub>p</sub> of <sup>208</sup>Pb determines P at low densities of about 2/3ρ<sub>0</sub> (average of surface and interior ρ).
- Radius of (~1.4M<sub>sun</sub>) NS depends on P at medium densities of ~2ρ<sub>0</sub>.
- Maximum mass of NS depends on P at high densities.



• These three measurements constrain density dependence of EOS and possible pressure changes from phase transitions.

#### PREX uses Parity Violation to measure $R_n(^{208}Pb)$

- In Standard Model Z<sup>0</sup> boson couples to the weak charge.
- Proton weak charge is small:  $Q_W^p = 1 - 4 \sin^2 \Theta_W \approx 0.05$
- Neutron weak charge is big:

 $Q_W^n = -1$ 

- Weak interactions, at low Q<sup>2</sup>, probe neutrons.
- Parity violating asymmetry A<sub>pv</sub> is cross section difference for positive and negative helicity electrons

$$A_{pv} = \frac{d\sigma/d\Omega_{+} - d\sigma/d\Omega_{-}}{d\sigma/d\Omega_{+} + d\sigma/d\Omega_{-}}$$

 A<sub>pv</sub> from interference of photon and Z<sup>0</sup> exchange. In Born approximation

$$A_{pv} = \frac{G_F Q^2}{2\pi\alpha\sqrt{2}} \frac{F_W(Q^2)}{F_{\rm ch}(Q^2)}$$

$$F_W(Q^2) = \int d^3r \frac{\sin(Qr)}{Qr} \rho_W(r)$$

 Model independently map out distribution of weak charge in a nucleus.

#### • Electroweak reaction free from most strong interaction uncertainties.

Donnelly, Dubach, Sick first suggested PV to measure neutrons.



PREX measures how much neutrons stick out past protons (neutron skin).

#### First PREX result and future plans

- At Jefferson Laboratory, I.05 GeV electrons elastically scattered from thick 208Pb foil. PRL 108, 112502, PRC 85, 032501
- A<sub>PV</sub>=0.66 ±0.06(stat) ±0.014(sym) ppm
- Neutron skin thickness:  $R_n-R_p=0.33^{+0.16}-0.18$  fm
- Experiment achieved systematic error goals.
- •Future plans: **PREX-II** (approved 25 days) Run <sup>208</sup>Pb again to accumulate more statistics. Goal:  $R_n$  to ±0.06 fm.
- •**CREX**: Approved follow on for <sup>48</sup>Ca with goal: R<sub>n</sub> to ±0.02 fm.



## Model dependent determinations of neutron skins

- Neutron skin thickness closely related to L (density dependance of symmetry energy).
- Light blue region in Lattimer et al from Japanese model dependent analysis of proton elastic scattering from Sn isotopes. Very old relativistic impulse approximation amplitudes from me were fudged so that <sup>58</sup>Ni had zero skin. Then amplitudes were used to describe Sn scattering. This is definitely model dependent!
- Danielewicz also has very precise Isobaric Analog State result that assumes same Sn skin analysis!





Gandolfi et al. PRC85, 032801 (2012)

# Fate of NS merger central object

- Prompt collapse to black hole, delayed collapse, or absolutely stable.
- Is very important EOS observable!
- Observation of black hole formation (with some knowledge of total system mass) sets UPPER limit on maximum NS mass, and rules out all very stiff EOSs.

#### SN neutrinos and r-process nucleosynthesis

- Half of heavy elements (including gold) are believed made in the rapid neutron capture process. Seed nuclei captures many n and decay. What makes all the neutrons?
- **Neutrinos:** Important possible site for the r-process is the neutrino driven wind in core collapse supernovae.
- Ratio of neutrons to protons in wind set by capture rates that depend on neutrino / antineutrino energies.

 $\nu_e + n \to p + e \quad \bar{\nu}_e + p \to n + e^+$ 

 $\Delta E = \langle E(\bar{\nu}_e) \rangle - \langle E(\nu_e) \rangle$ 

 Measure ΔE, difference in average energy for antineutrinos and neutrinos. If ΔE is large, wind will be neutron rich. If ΔE is small, wind will be proton rich and likely a problem for r-process.



Searching for El Dorado with supernova neutrinos

Important to measure energy of both anti-nu (SK) and neutrinos (liquid argon?).

 $\Delta E$  depends on some nuclear physics including symmetry E at low densities.

#### However, present SN simulations find too few neutrons.

### Supernova Neutrino Detectors

Are crucial to our whole field!

Detector	Type	Mass (kt)	Location	Events	Live period	
Baksan	$C_n H_{2n}$	0.33	Caucasus	50	1980-present	
LVD	$C_n H_{2n}$	1	Italy	300	1992-present	
Super-Kamiokande	$H_2O$	32	Japan	7,000	1996-present $\bar{\nu}_e + p \rightarrow$	$e^+ + n$
KamLAND	$C_nH_{2n}$	1	Japan	300	2002-present	
MiniBooNE*	$C_n H_{2n}$	0.7	USA	200	2002-present	
Borexino	$C_n H_{2n}$	0.3	Italy	100	2005-present	
IceCube	Long string	0.6/PMT	South Pole	N/A	2007-present	
Icarus	Ar	0.6	Italy	60	Near future	
HALO	Pb	0.08	Canada	30	Near future	
SNO+	$C_n H_{2n}$	0.8	Canada	300	Near future	
MicroBooNE*	Ar	0.17	USA	17	Near future	
$NO\nu A^*$	$C_nH_{2n}$	15	USA	4,000	Near future	
LBNE liquid argon	Ar	34	USA	3,000	Future $\nu_e + {}^{40}\text{Ar} \rightarrow$	$e^{-} + {}^{40}\text{K}^{*}$
LBNE water Cherenkov	$H_2O$	200	USA	44,000	Proposed	
MEMPHYS	$H_2O$	440	Europe	88,000	Future	
Hyper-Kamiokande	$H_2O$	540	Japan	110,000	Future	
LENA	$C_nH_{2n}$	50	Europe	15,000	Future	
GLACIER	Ar	100	Europe	9,000	Future	

Supernova neutrino detectors. Events from a SN at 10kpc (considerable model variation). — Kate Scholberg

Important to measure spectrum of all three components: electron neutrinos, electron antineutrinos, and nu-x. Can measure nu-x in elastic scattering detector.

#### Long Baseline Neutrino Experiment



- Send neutrino and antineutrino beams from Fermilab (near Chicago) 1300 km to a large (34 kt) liquid Ar detector in the Homestake gold mine in South Dakota. Main goal: observe CP violation in neutrino oscillations.
- Powerful supernova detector that should be able to measure electron neutrino energies very well.
- Combine anti-nu E from Super K with nu E from LBNE to predict composition of neutrino driven wind and likely strongly disfavor rprocess in wind.

#### Detecting extra-galactic supernova neutrinos in the Antarctic ice (10 Megaton detector)

- Preliminary idea to instrument inner region of IceCube to obtain 10 MeV threshold and 10 Mt effective volume.
- Issues with cost, photodetectors, and noise.
- See SN to 10 Mpc with rate of 1 to 4 per year.
- Boser et al, arXiv:1304.2553
- Coincident with GW signal.
- What do you learn from handful of events?
- What else can this detector do?



#### SN neutrinos and r-process nucleosynthesis

 Ratio of neutrons to protons in wind set by capture rates that depend on neutrino and anti-neutrino energies.

 $\nu_e + n \to p + e$   $\bar{\nu}_e + p \to n + e^+$ 

- Weak magnetism important
- Measure difference in average energy of antineutrinos and neutrinos. If large, wind will be neutron rich. If it is small, wind will be proton rich and likely a problem for r-process.
- Composition (Y<sub>e</sub>) of wind depends on anti-neutrino energy (Y-axis) [results from ~20 SNI987A events shown] and energy of neutrinos (X-axis). Energy of neutrinos, not yet measured, depends on properties of n rich gas (nu-sphere).



LBNE measures X axis and Super-K measures Y axis.

Present SN simulations find too few neutrons for (main or 3rd peak) r-process. Suggests this is not r-process site. However composition still determines what SN wind makes!

### Neutrinosphere Problem

• Build detailed neutrino atmosphere model with quantifiable uncertainties to predict  $\nu_{e}$ , anti- $\nu_{e}$ , and  $\nu_{x}$  spectra for:

- SN  $\boldsymbol{\nu}$  detectors

- Input to  $\boldsymbol{\nu}$  oscillation calculations
- Electron fraction  $Y_e$  of  $\boldsymbol{\nu}$  driven wind and nucleosynthesis
- Focus on  $\Delta E = \langle E(anti-\boldsymbol{\nu}_e) \rangle \langle E(\boldsymbol{\nu}_e) \rangle$

### "Femtonovae"

- Core collapse SN dominated by neutrinos. Much of the "action" occurs near the neutrinosphere (surface of last scattering) at temperatures of ~5 MeV, densities of 1/1000 to 1/10 ρ<sub>0</sub>, and neutron rich compositions.
- A "Femtonova" is a very small new star. Suggested name for HI collisions applied to astrophysics, and in particular to recreate neutrinosphere conditions.
  - Study in the laboratory the equation of state, symmetry energy, composition, and neutrino response ... of neutrinosphere material.
    - Recreating ~5 MeV temperature is straight forward.

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- Recreating low densities occurs as system expands but it may be difficult to measure the density.
- Recreating the very neutron rich conditions is harder. Perform HI collisions with proton rich and then neutron rich radioactive beams and extrapolate to very neutron rich conditions.

### Chemical Freeze Out

- HI collisions create warm source regions that then expand with time.
- Nucleons in source regions interact to form light clusters D, <sup>3</sup>H, <sup>3</sup>He, <sup>4</sup>He.
- Source expands to such low densities that clusters stop interacting and the chemical composition freezes out.
- One can then measure this final composition which should correspond to the composition of equilibrium nuclear matter at the freeze out temperature and density.
- Several ways to measure freeze out temperatures (often near 5 MeV). For example looking at ratio of yields.  $T_{HHe}$ =14.3/ ln[1.59 (Y<sub>D</sub>Y<sub>4He</sub>/Y<sub>3H</sub>Y<sub>3H</sub>e)]
- Measuring freeze out densities are more difficult. Example, can get source sizes from two-particle correlation functions.

#### Virial expansion for Neutrinosphere EOS

- Start with pure neutron matter, then include clusters.
- Assume: system in gas phase and has not undergone a phase transition with increasing density or decreasing T and that the fugacity z=e<sup>µ/T</sup>with µ the chemical pot is small.
- Expand pressure in powers of z :

 $P=2T/\lambda^{3}[z+b_{2}z^{2}+b_{3}z^{3}+...],$ 

Here  $\lambda$ =thermal wavelength=(2/mT)<sup>1/2</sup>

• 2<sup>nd</sup> virial coef. b<sub>2</sub>(T) from 2 particle partition function which depends on density of states determined from phase shifts:

$$b_2 = 2^{1/2} \sum_B e^{E_B/T} + \frac{2^{1/2}}{\pi} \int_0^\infty dk \, e^{-E_k/2T} \sum_l (2l+1) d\delta_l(k) / dk \pm 2^{-5/2}$$

### Nuclear Matter: *n*, *p*, $\alpha$ system

$$\frac{P}{T} = \frac{2}{\lambda^3} [z_p + z_n + (z_n^2 + z_p^2)b_n + 2z_n z_p(b_{nuc} - b_n)] + \frac{1}{\lambda^3_{\alpha}} [z_{\alpha} + z_{\alpha}^2 b_{\alpha} + z_{\alpha}(z_p + z_n)b_{\alpha n}]$$

- Need four virial coefficients:
  - b<sub>n</sub> for neutron matter,
  - b<sub>nuc</sub> for symmetric nuclear matter,
  - $b_{\alpha}$  for alpha system,
  - $b_{\alpha n}$  for interaction between an  $\alpha$  and N.
- Virials from NN, N $\alpha$  and  $\alpha \alpha$  elastic scattering phase shifts.



 $\alpha$ - $\alpha$  Elastic Phase Shifts

#### Recreating Neutrinosphere on Earth

- Can study neutrinosphere like conditions with heavy ion collisions in lab.
- Composition of intermediate velocity fragments in 35 MeV/n <sup>64</sup>Zn on <sup>92</sup>Mo and <sup>197</sup>Au: Data (blue squares) Kowalski et al, PRC 75, 014601(2007). Virial EOS black

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Target



In a peripheral HI collision, intermediate velocity fragments from warm low density region.

Intermediate velocity fragments

### Symmetry Energy shift

- Proton in n rich matter more bound than neutron because of symmetry energy.
- Symmetry energy at low density can be calculated exactly with virial expansion (with A. Schwenk). Find it is much larger than in some mean field models because of cluster formation.
- Neutrino absorption cross section increased by energy shift which increases energy and phase space of outgoing electron-> lowers E(nu).
- Consider  $\nu_e$  + n -> p + e

$$\begin{aligned} \Delta U &= U_n - U_p = \lambda^3 T (n_n - n_p) (b_{pn} - \hat{b}_n) \\ \\ \frac{\sigma_{\nu_e}(\Delta U)}{\sigma_{\nu_e}(0)} &= \frac{(E_\nu + \Delta U)^2 [1 - f(E_\nu + \Delta U)]}{E_\nu^2 [1 - f(E_\nu)]} \end{aligned}$$

- Effect opposite for anti-neutrino absorption and reduces cross section increasing E(anti-nu).
- Increases ΔE and makes wind somewhat more neutron rich. Probably not enough for r-process ?? But symmetry energy is relevant.

Idea due to L. Roberts, my work with G. Shen, C. Ott, E. O'Connor



Symmetry energy from isoscaling analysis of ratio of yields of light clusters with different N/Z values. The temperature varies from about 4 MeV (lowest density) to 10 MeV (highest density)

#### EOS and Neutrino Interactions for Nucleosynthesys

- Combine information from Super-K for anti- $\nu_e$  and LBNE for  $\nu_e$  from a galactic supernova, allows one to infer the composition of nu driven wind providing more accurate nucleosynthesis predictions.
- Should make detailed neutrino atmosphere model to predict  $\nu_{e}$ , anti- $\nu_{e}$ , and  $\nu_{x}$  spectra for: SN  $\nu$  detection,  $\nu$ -osc calculations,  $\nu$ -wind Y<sub>e</sub> and nucleosyntheses.
- Femtonovae: Can recreate neutrinosphere conditions in lab with HI collisions and measure EOS, composition, symmetry E, nu-response to benchmark nu-atmosphere model.
- Both HI collisions and Virial EOS find large Symmetry E (or self E) shift that will increase  $\nu_e$  and reduce anti- $\nu_e$  charged current cross sections.
- Collaborators: G. Shen, E. O'Connor, C. Ott, A. Schwenk, Z. Lin... Supported in part by DOE grants DE-FG02-87ER40365 (Indiana U.) and DE-SC0008808 (NUCLEI SciDAC).

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